PA - 2126 Investigation of the penetrating plasma on the occasion of a discharge in mercury vapors.

determined only by the amount of the wall flux of the ions and is independent of the amount of the negative potential.

ASSOCIATION:

Physical Institute of the Academy of Science of the U.S.S.R., Kiev

PRESENTED BY:

SUBMITTED: 6.1956

AVXILABLE:

Library of Congress.

Card 3/3

E GABOVICH, M. D.

"Anomalous Scattering and Excitations of Plasma Oscillations and Plasma Resonance."

paper presented at Second All-Union Conference on Gaseous Electronics, Moscow, 2-6 Oct '58.

TO SECURE DESCRIPTIONS DEPOSITE AND ACTION OF THE PROPERTY OF

OABOVICH, M.D. [Gabovych, M.D.]; NEMETS, O.F.; FEDORUS, Z.P.

On the utilization of a high-current pulse discharge in proton sources [In Ukrainian with summary in English]. Ukr.fiz.zhur. 3 no.1:104-111 Ja-7 '58. (MIRA 11:4)

1.Institut fiziki AN URSR. (Protons) (Electric discharges through gases)

Critical effect of the geometric dimensions of a high-fraquency proton source system on the current intensity of an ion beam.

Ukr. fiz. zhur. 3 no.3:419-421 My-Je '58. (MIRA 11:10)

1. Institut fiziki AN USSR. (Protons) (Ion beams)

57-28-4-31/39 Gabovich, M. D. AUTHOR ! On the Rational Utilization of the Magnetic Field in High-PITLE -Frequency Sources of Protons (O ratsional nom ispol'zovanii magnitnogo polya v vysokochastotnykh istochnikakh protonov) Zhurnal Tekhnicheskoy Fiziki, 1958, Vol. 28, Nr 4, pp. 372-PERIODICAL: -880 (USSR) The influence of the magnetic field upon a high-frequency ABSTRACT: discharge is investigated here and in this connection some conclusions on the rational utilization of the magnetic field in high-frequency sources of protons for practical purposes are drawn. Of the three possible kinds of utilization of the magnetic field in high-frequency sources that of the so--called "transverse effect" of the magnetic field (references 1 3) is investigated. All facts given here refer to the frequencies of 40-50 megacycles. It is shown that during the influence of the transverse magnetic field upon a high--frequency discharge, which is excited by means of an inductive connection, this field does not only promote the introduction of the power into the domain of discharge, but also leads to a much more effective utilization of this Card 1/2

On the Rational Utilization of the Magnetic Field in High- 57 28-4-31/39

power. It is further shown that this relatively weak magnetic field at the same time brings about an anisotropy in the distribution of the density of the ionization current. The latter fact leads to the conclusions that the mutual position of the coil causing the discharge, of the magnetic field and of the axis from which the ions are drawn out could be more rational. This conclusion is confirmed by the experiments with an ion source. It is shown that the use of a heterogeneous end-magnetic-field is more expedient than the use of a homogeneous magnetic field. Some experimental data on the influence of the magnetic field upon the high-frequency discharge are given. - Yu.S. Alpat'yev and Z.P. Fedorus participated in the measurements. There are 8 figures, 2 tables, and 7 references, 2 of which are Soviet.

ASSOCIATION:

Institut fiziki AN USSR, Kiyev (Kiyev, Institute for Physics,

AS Ukrainian SSR)

SUB TTTED:

December 14, 1956

Card 2/2

AUTHORS: Granovski, V.I. AUTHORS: Granovski, V.I. FIRITONICAL: Redictionals AMAINISTY The conference AMAINISTY OF MAR. A		Grandwakiy, V.L., Luk.yanov, S.fu., Spivek, G.V. and Sirecako, X.G. Sirecako, X.G. Raport es the Second All-Union Conference on Gas	Hadio	H. P. GADENTIAN H. P. GADENTIAN TALE, TILEMEN THA "EXCITATION TO "A	The section of the Stark Education of Stark Educ	
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SOV/56-36-4-10/70 21(7) Gabovich, M. D., Pasechnik, L. L. AUTHORS: The Anomalous Scattering of Electrons and the Excitation of TITLE: Plasma Oscillations (Anomal'noye rasseyaniye elektronov i vozbuzhdeniye plazmennykh kolebaniy) Zhurnal eksperimental noy i teoreticheskoy fiziki, 1959, PERIODICAL: Vol 36, Nr 4, pp 1025-1033 (USSR) In the introduction, several papers dealing with this subject ABSTRACT: are discussed (Refs 1-11). The object of the present paper was the investigation of interaction between the electron beam and a plasma formed independently. The experimental arrangement (Fig 1) consisted essentially of a glass tube and an attached piece containing a probe. By means of a liquid mercury cathode and a special anode system a plasma was produced along the tube in the mercury vapor, the density of which amounted to $1.10^9 - 1.5.10^{11}$ cm⁻³ (mercury vapor pressure p $\approx 1.10^{-3}$ torr). An oxide cathode served as electron source. First, the characteristic at various (small) currents I of the electron beam Card 1/4 was investigated; figure 2 shows the dependence of the collector

SOV/56-36-4-10/70
The Anomalous Scattering of Electrons and the Excitation of Plasma Oscillations

current on the grid voltage on the analyzer probe (8 curves for I values of 0.05-13 ma with an anode current of 0.5 a, E=50v, l=50 mm). Plasma concentration was $n=1.6\cdot10^{10}$. The conclusions drawn from the course of the curves and the phenomena of anomalous scattering are discussed (The phenomenon which was first observed by Langmuir (Ref 1), consists in principle in the fact that electrons which have penetrated the plasma partly have high velocities). Figure 3 shows the dependence of the relative quantity of anomalously fast electrons on the position of the probe (again for different I_e -values).

The problem of the limiting current is discussed and illustrated by a table for different types of cathodes and different ancel currents and plasma densities \mathbf{I}_{lim} and \mathbf{j}_{lim} . The following chapter discusses excitation and extinguishing of plasma cscillations occurring as a results of interaction between the electron beam and the plasma. Figure 4 shows the spatial course of oscillation intensities at various \mathbf{I}_{e} -values and constant

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 $I_a = 10 \text{ ma}$, E = 40v. Figure 5 shows oscillation intensity

sov/56-36-4-10/70

The Anomalous Scattering of Electrons and the Excitation of Plasma Oscillations

distribution at various I -values and constant I = 13.5 ma. Figure 6 shows the dependence of wave length and oscillation intensity on $I_{\rm R}$ and figure 7 finally shows the radial intensity distribution of oscillations in the electron beam penetrating the plasma for various I, -values. In the following, the influence exercised by an external magnetic field oriented parallel to the electron beam is discussed in short. Figure 8 shows the dependence of the position of the oscillation zone and of the scattering zone on the wave length of the observed electromagnetic oscillations (1 increases linearly with A); figure 9 shows the dependence of the position of the scattering zone on λat 2 electron energies, E = 41 and 28 v. 1 also grows linearly with A, the curve for greater E is somewhat steeper. Figure 10 shows the same for λ -values corresponding to a certain plasma concentration. A discussion of the results obtained shows that the effects observed may be explained qualitatively by the fact that electrons are assumed to gather in clusters and that these clusters coherently interact with the plasma. The

Card 3/4

SOV/56-36-4-10/70

The Anomalous Scattering of Electrons and the Excitation of Plasma Oscillations

authors finally thank N. D. Morgulis for discussing results. There are 10 figures, 1 table, and 15 references, 4 of which

are Soviet.

ASSOCIATION:

Institut fiziki Akademii nauk Ukrainskoy SSR (Physics Institute

of the Academy of Sciences, Ukrainskaya SSR)

SUBMITTED:

October 10, 1958

Card 4/4

9.3150,24.2120

TITLE:

77847 \$<mark>0</mark>707/57-30-3-13/15

AUTHORS: Gabovich, M. D.,

Gabovich, M. D., Bartnovskiy, O. A., Fedorus, Z. P.

Sag of the Potential on the Axis of a Discharge at

Electron Oscillation in a Magnetic Field

PERIODICAL: Zhurnal tekhnicheskoy fiziki, 1960, Vol 30, Nr 3,

pp 345-350 (USSR)

ABSTRACT: Kistemaker and Sheider (Physica, 19, 950, 1953) showed

experimentally that in a discharge with electron oscillations in magnetic field potential on the axis of discharge may be considerably smaller than potential of anode. In the present paper the authors investigate causes for such a potential sag and examine conditions favoring effect. Figure 1 shows the diagram of experimental setup and measuring circuitry. In addition to cathode K and anode A, there are two reflectors 0 and 0 at the potential of the cathode of negative

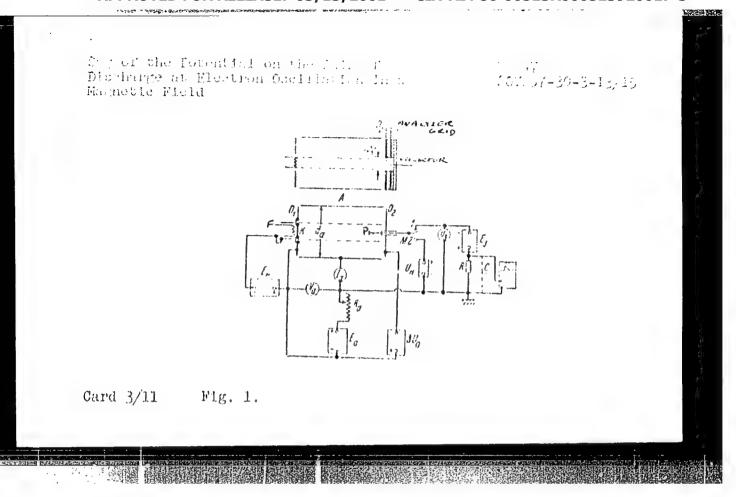
with respect to it. The cathode was either of tantalum, indirectly heated by bombardment of electrons originating on F or a directly-heated tungsten cathode. The

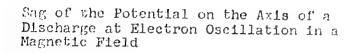
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77847 \$0V/57-30-3-13/15

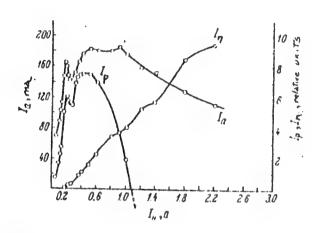
whole 35 mm length of the system was in an uniform longitudinal magnetic field H variable 0-4,000 oersted. The behavior of anode current \mathbf{I}_a , probe current \mathbf{I}_p (at -80 v with respect to anode) and noise intensity in probe curcuit \mathbf{I}_n as functions of magnetic field are presented in Fig. 2. For $\mathbf{I}_H=1$, H ≈ 500 oersted. U a was 300 v with respect to the cathode. The authors prove irregularities of the \mathbf{I}_a curve are unambiguously related to noise intensity. They explain these irregularities by formation of a fundamental discharge column caused by axial oscillations of primary electrons in the raising magnetic field. At a certain optimum value of \mathbf{I}_H the field starts substantially preventing plasma electrons from reaching the anode and produces a potential "groove." Its radial electrical field, in turn, facilitates motion of electrons toward the anode which was hampered by the presence of

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Card 4/11

Fig. 2.

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magnetic field. Further increase of H produces an unstable discharge, causing the mentioned irregularities and noises. The probe current changes sign because of an increasing number o' primary electrons reaching it and a decrease of potential or paraxial plasma. Further increase of the magnetic field finally takes over and decreases the anode current until discharge is apparently completely halted. To measure potential inside the plasma the authors developed a special thermal probe consisting of a tungsten disc 1 mm diam and 0.05 mm thick on a tungsten wire inside an insulating quartz tube. By a relay M (see Fig. 1) probe P is raised to a potential V_H during a time interval $\mathcal T$ The electron current bombarding the probe can heat it sufficiently to produce an appreciable electron emission. During the second half of the cycle \mathcal{T}_2 probe is at potential Up and, if the heating effect is now lower than previously, emission will decrease. Now, in the

Card 5/11

Sag of the Potential on the Axis of a Discharge at Electron Oscillation in a Magnetic Field

77847 SOV/57-30-3-13/15

case of U_p being lower than plasma potential, decrease of emission is accompanied by a decrease of probe current while in $U_p \supset U$ plasma current changes sign and remains constant in time. The authors changed probe potential 20 times per second, observed current pattern on an oscilloscope, and registered plasma potential from those readings of the U_p voltmeter at which the decaying current pattern on the oscilloscope screen switched to the rectangular one. Results for measured potential U_a and plasma potential on discharge axis U_a are shown in Fig. 6 as a function of magnetization current I_H and diam of the anode. Analysis of results showed $\Delta U_a = U_a - U_b$ is a linear function of the square of the anode diam:

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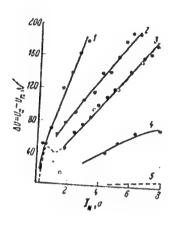


Fig. 6. (1) $d_a = 4.0 \text{ cm}$ (2) $d_a = 3.4 \text{ cm}$ (3) $d_a = 2.7 \text{ cm}$ (4) $d_a = 1.8 \text{ cm}$ (5) $d_a = 1.0 \text{ cm}$

Card 7/11

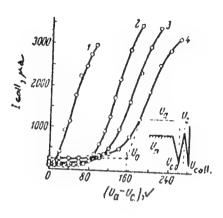
77847 307/57-30-3-13/15

The authors discovered that radial potential drop is almost completely located outside the axial plasma of diam equal to diam of the cathode. They note, however, all measurements mentioned above were done in the presence of a perturbation caused by the presence of the probe. They circumvent this objectionable situation by developing a special setup consisting of a grid across an $\phi = 8$ mm opening on the reflector 0_{2}

followed by another analyzer grid and a collector. Distribution of potentials is shown on the right in Fig. 8. The authors assumed there would be an apprecable ion current on the collector only when potential of analyzer grid U_c is equal or smaller than potential of plasma U_n. Using these values they constructed the curves in Fig. 8 for an anode 2.7 cm diam. Extrapolated potential values in the manner indicated in Fig. 8 then yielded points marked by hollow circles in Fig. 6. The agreement between the two methods is apparently very good.

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77847 \$0V/57-30-3-13/15



Card 9/11

Fig. 8. (1)
$$I_H = 1.5 \text{ a}$$
 (2) $I_H = 3.5 \text{ a}$ (3) $I_H = 5.0 \text{ a}$

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Sag of the Potential on the Azis of a Discharge at Electron Oscillation in a Magnetic Field

77847 SOV/57-30-3-13/15

The authors finally did some theoretical calculations starting from the equation of radial electron current density

$$j_{-} = -D \frac{dn_{-}}{dr} + \frac{Dc}{kT} n_{-} \frac{dU}{dr}$$
 (1)

and the continuity equation

$$\frac{dj_{r-1}}{d\tilde{r}} \stackrel{j_{r-1}}{=} \beta n_{r-1} \tag{2}$$

Assuming n_ to be constant, they obtained a theoretical expression for $\Delta\, U$ in volts

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Sag of the Potential on the Axis of a Discharge at Electron Oscillation in a Magnetic Field

77847 \$0V/57-30-3-13/15

 $\Delta U(\mathbf{v}) \approx 10^{-2} \cdot H(\mathbf{c}) d_a^2$

(5)

which for H = 1,500 oersted and d_a = 4 cm yields \triangle U \cong 240 v versus the experimentally measured value 180 v. The authors note relationship U = f(H, d_a) as well as value \triangle U are in fair agreement with the experiment. The strong radial fields up to 100 v/cm are connected to a decrease of electron diffusion towards the anode. There are 8 figures; and 6 references, 3 Soviet, 1 Dutch, 1 German, and 1 U.S. The U.S. reference is: D. Bohm. The Characteristics of Electrical Discharges in Magnetic Fields. N. Y. 1949.

ASSOCIATION:

None given

SUBMITTED:

April 18, 1959

Card 11/11

9.3150,21.2100,24.2120

77848 SOV/57-30-3-14/15

AUTHOR:

Gabovich, M. D.

TITLE:

Role of Multiple Processes in Proton Formation in Ion

PERIODICAL:

Zhurnal tekhnicheskoy fiziki, 1960, Vol 30, Nr 3, pp

354-358 (USSR)

ABSTRACT:

At the present time it is possible to obtain ion beams with a large proton content from appropriately constructed ion sources. The proton production mechanism 1s, however, in most cases unexplained except that experimental data indicate direct (single) processes during electron-molecule collision cannot represent an effective mechanism of proton creation. The author assumes a large proton content is connected to multiple processes of electron collisions with heavy particles and considers some of the possibilities together with experimental data in favor of such a point of view.

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He investigates the following possibilities:

a)
$$H_2$$
-1· $e \rightarrow H_1$ -1· H_0 b) H_1 -1· $e \mapsto H_1^+$

a)
$$H_2 + e \rightarrow H_2^+$$
, b) $H_2^+ + e \rightarrow H_1^+ + H_1^-$

a)
$$H_2 \rightarrow e \rightarrow H_3^+$$
, b) $H_3^+ \rightarrow e \rightarrow H_1 \rightarrow H_1$, c) $H_1 \rightarrow e \rightarrow H_1^+$.

The author assumes also: (1) gas pressure and degree of ionization are low, particle free paths are larger than size of the container, and gradients of particle concentrations are negligibly small; (2) recombination of atomic hydrogen occurs only on the walls and represents an effect of the first order; the number of recombined atoms is proportional to the product of the flow of atoms toward the walls and the coefficient of recombination. The ratios of the numbers of obtained protons and molecular ions are then given by:

for process 1

$$\frac{V_1 + = 0 \text{ in Final}}{3 \text{ var}} \int_{\text{Fidel}} h_{\text{Ha}}$$

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for process 2
$$\frac{V_{1:1}}{V_{2:1}} = \frac{n_{H_1}}{n_{H_1}} \cdot \frac{\int_{diss_{H_1}}}{\int_{j\partial n_{H_1}}}$$
 (2)

for process 3
$$\frac{v_{1+}}{v_{2+}} = \frac{8RFn_s}{3v_{ar}} \cdot \frac{n_{H^+}}{n_{H_1}} \int_{dist}$$
 (3)

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Here $n_{\mbox{\scriptsize e}},~n_{\mbox{\scriptsize H2}}$ and $n_{\mbox{\scriptsize H2}}$ are respective concentration of

77848 sov/57-30-3-14/15

e, H_2^{\dagger} , and H_2 ; \mathcal{V}_2 is mean velocity of atoms. \int diss. H_2 represents the integral

 $\int \sigma dss_{H_1} c f(c) dc,$

with diss. H₂ as electron collision dissociation cross-section of molecules, and f(c) as velocity distribution function for electrons. F represents ratio of integrals corresponding to atomic and molecular ionization. In low ionization, as in the plasma of a stationary hf ion source with moderate generator power, the dominant production mode is (1). The author computes the process using the diss. H₂ value by Massey

Card 4/9

and Mohr. (Proc. Roy. Soc., A, 135, 358, 1932), assumes

77548 30V/57-30**-3-14/1**5

f(c) to be Maxwellian, and obtains

 $B = \frac{\beta}{(1-\beta)(1-0.3\beta)} \frac{76dD(kT_s)F[^{*}[_{t_0}]^{*}]}{[_{t_0}](s_0/s_0c_s)(S_1r_1+c_0r_2)(s_0r_1)(s_0r_1)} \hat{J}_{+}[_{s_0}]^{*} (7)$

where /3 is given by

 $j_{i+1}, j_{i+2+j_{i+3}}$

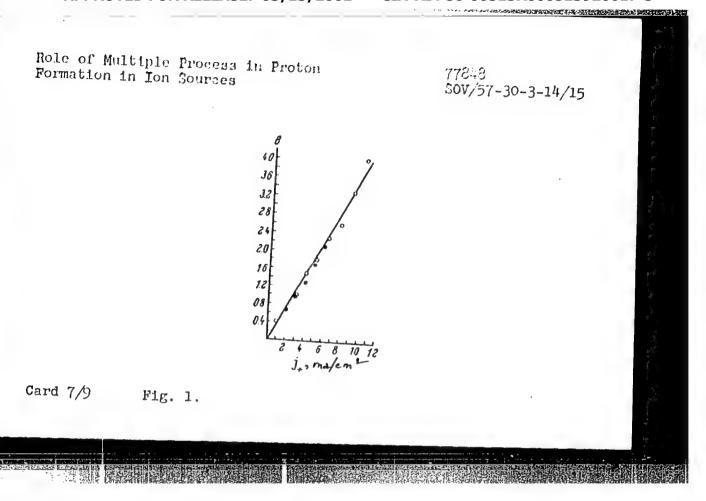
 J_{+k} is current densities corresponding to lons H_1^+ , H_2^+ , H_3^+ ; V is volume of the source; S_1 is its glass surface; S_2 is surface of metallic parts; r_1 and r_2 are corresponding coefficients of recombination; $G(kT_e)$ is certain tabulated functions slowly varying with kT_e in the interval under consideration (for $T_e = 5 \cdot 10^{\frac{1}{4}}$ - $1 \cdot 10^{\frac{1}{4}}$ or, $G(kT_e)$ varies from 0.5-0.62). According

Card 5/9

77848 80V/57-30-3-14/15

to Eq. (7) B should be proportional to the total ion current density. To check the expression, the author established a 45 Mc/s discharge of variable intensity inclde a spherical pyrex tube. Using two electrode probes (Al and Mo) and a mass analyzer described elsewhere (M. D. Gabovich. PTE, 2, 88, 1996), the author measured composition of the Ion beam and obtained the curve in Fig. 1 which is indeed linear. Performing auxiliary experiments with and without a tungsten spiral inside the discharge tube, the author established the ratio r_2/r_1 is in the limits 40 - r_2/r_1 80. Taking $r_1 = 1.6 \cdot 10^{-4}$ from the work of Poole (Proc. Roy. Soc., $1\overline{6}3$, 404, 1937) the author found r_2 to be significantly smaller than 1 so he could neglect the term $\mathbf{S}_{_{\mathcal{O}}}\mathbf{r}_{_{\mathcal{O}}}$ in Eq. (7). Computing then theoretically he got a value for $B/J_+ = 0.44$. Comparing this value with the experimental value of 0.37, he concluded the assumption about the paramount importance of the process (1) was right.

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77848 SOV/57-30-3-14/15

The author further concluded to obtain a value of B 80%, which was reached in the pyrex tube experiment with $J_{+} \approx 10 \text{ ma/cm}^2$, one would need 1-10 a/cm² if the tube were metal because of strong recombination on This agrees with experimental results of Gabovich and others (UFZ, 3, 104, 1958); although there, in view of a larger degree of ionization, competing processes could have played an important role. In the case of proton production by a metallic capillary arc, the proton content is limited by the small value of the effective radius of the chamber $R_{\text{eff}} \sim V/S$ in Eq. (7). Competing with the above electron collision processes in the proton production are collision processes between heavy particles. The widespread belief that the latter processes are characterized with a large cross-section are not experimentally verified in the electron-volt energy region. The author also notes precise mass-spectrometric results of Demirkhanov (M. Ardenne. Tabellen der Elektronenephysik, Jonenphysik u. Ubermikroskopie, (Tables of the Physics of

Card 8/9

love of Multiple Processes in Proton Formation in Ion Sources

77343 SOV/57-30-3-14/19

Flectrons, Ions, and Ultra-Microscopes) Berlin II, 836, 1956) who showed that in a plasmatron, along with fast protons caused by process (1), one finds slow protons because of the disintegration of Hg. However, in

the coefficient of recombination is smaller than 1, the slow protons could also be caused by atoms that, by the time they become ionized, lose most of dissociation energy in collision with the walls. There are 2 figures; and 16 references, 4 Soviet, 1 German, 5 U.K. and 6 U.S. The recent U.K. and U.S. references are: B. J. Wood, H. Wise. J. Chem. Phys., 29, 1416, 1958. L. A. Edelstein. Nature, 182, 932, 1958. E. V. Ivash. Pnys. Rev., 112, 155, 1958. E. Bauer, Ta-You-Wu. Canad J. Res., 34, 1436, 1956. E. H. Kerner. Pnys. Rev., 92, 1441, 1953.

ASSECTATION: Institute of Physics AS UkrSR, Kiyev (Institut fiziki

AN USSR, Kiyev)

SUBMITTED.

March 8, 1957

Cara 9/9

"APPROVED FOR RELEASE: 03/13/2001

CIA-RDP86-00513R000513920017-3

83577

9,6/5C 26.2310 24.2120

S/056/60/038/005/010/050 B006/B070

AUTHORS:

Gabovich, M. D., Pasechnik, L. L., Yazeva, V. G.

TITLE:

Detection of Ion Oscillations in a Plasma ?

PERIODICAL:

Zhurnal eksperimental noy i teoreticheskoy fiziki, 1960,

Vol. 38, No. 5, pp. 1430-1433

TEXT: Ion oscillations with a limiting frequency of $f_0 = \sqrt{ne^2/\pi M}$ have been known for electron beams with compensated space charge, but they had not yet been found in the plasma of a gas discharge. It is shown in the present work that it is possible to make a direct determination of self-sustaining ion oscillations in the plasma of a gas discharge. The experimental apparatus consists of a discharge tube in which there is an arc discharge in mercury vapor; the charge concentration in the plasma can be varied by varying the discharge current. There are two probes in the plasma, one fixed and the other movable. The distance between them could be altered from 0 to 15 mm. The arrangement for the detection of ion oscillations is described in brief. Essentially, it consists of a preamplifier, a superheterodyne amplifier of the type $M\Pi-12M$ (IP-12M), a special three stage Card 1/3

83577

Detection of Ion Oscillations in a Plasma

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narrow-band amplifier, and a tube voltmeter. The sensitivity of the amplifying arrangement can reach $\sim 2\cdot 10^{-8}$ v. The results of measurement are shown in Fig. 2: With increasing discharge current I, the voltage U_{out} at the output of the amplifier system increases, passes through a maximum, and then falls steeply. The position and the height of the signal peaks in the $U_{out}(I)$ diagram depend on the frequency f of the amplifier. Fig. 2 shows the characteristics for f = 1.6, 2, and 2.4 Mc/sec. Fig. 3 shows the dependence of the resonance currents on the potential of the probes for f-values between 1.6 and 2.6 Mc/sec. I res increases linearly with U_{out} and the greater f the greater is the slope of this straight line. (I is the I-value corresponding to the peak of U_{out}). The following relation (2) holds for the frequency of the ion oscillations: $f = f_{o}/\sqrt{1 + ne^2 \lambda^2/\pi kT_e}$, where λ is the wavelength. With this, the charge density in the plasma is $V_{out} = V_{out} = V_{ou$

83577

Detection of Ion Oscillations in a Plasma

S/056/60/038/005/010/050 B006/B070

approximately equal to the radius of the ion layer surrounding the probe. Since the radius of the ion layer surrounding the probe increases with increasing potential of the probe, n and I must increase not only with f but also with negative potential U probe of the probe. This is actually found to be so experimentally. It is also found that $\lambda^2 < kT_e/Mf^2$. As a practical example (corresponding to the experimental conditions), one has $\lambda_{max} = 6.4 \cdot 10^{-2}$ cm with $T_e = 3.8 \cdot 10^4$ ok and $f = 2 \cdot 10^6$ cps. Such a thickness of the ion layer ($\sim \lambda_{max}$) fairly agrees with the experimental results. By extrapolating the curves shown in Fig. 3 for a zero potential of the probe, n and I ores may be obtained; and also here theory and experiment agree satisfactorily (Fig. 4). It has, thus, been possible to detect by these experiments the oscillations of ions and to verify formula (2) qualitatively. V. D. Rutgayzer and K. I. Kononenko are mentioned. There are 4 figures and 6 references: 1 Soviet, 4 US, and 1 Irish.

SUBMITTED:

November 23, 1959

Card 3/3

"APPROVED FOR RELEASE: 03/13/2001

CIA-RDP86-00513R000513920017-3

9,2585 9.4220

21/1/1/1

2/11 3/61/006/001/023/023 E140/E163

AUTHOR:

Gabovich, M.D.

TITLE:

On the mechanism of exciting electron plasma

oscillations

PERIODICAL: Radiotekhnika i elektronika, Vol.6, No.1, 1961,

pp. 178-179

In this letter the author discusses the conflicting TEXT: opinions concerning the significance of Merill and Webb's (Ref.1) fundamental experiments. The author expresses the opinion that only the klystron model of excitation of oscillations is capable of explaining all the known facts.

There are 23 references: 11 Soviet and 12 English.

ASSOCIATION:

Institut fiziki, AN USSR

(Physics Institute, AS Ukr.SSR)

SUBMITTED:

April 21, 1960

Card 1/1

GABDVICH, M.D.

27960 S/185/61/006/004/002/015 D274/D303

26,2331 AUTHORS:

Gabovych, M.D. and Romanyuk, L.I.

TITLE:

Plasma ejection by electrodeless pulse-discharge

in a vacuum

PERIODICAL:

Ukrayins'kyy fizychnyy zhurnal, v. 6, no. 4, 1961,

461-466

TEXT: Experiments are described with bursts of plasma in a narrow channel. Fig. 1 shows the experimental setup used. In glass tube 1 a discharge takes place at a pressure of $4 \cdot 10^{-2}$ mailg in region 1, and of below $1 \cdot 10^{-4}$ mmlig in region 5. The current related to the plasma-bursts was measured by moving electrode 6. Oscillograms were taken of the voltage drop across resistor 3 (which is part of the same circuit as 6). Oscillograms are shown of the current in circuit 6 at various voltages of the capacitor battery. The same figures show a plot of the time-derivatives of the exciting field-strength H. Another figure shows the dependence of the electron

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27960 S/185/61/006/004/002/015 D274/D303

Plasma ejection...

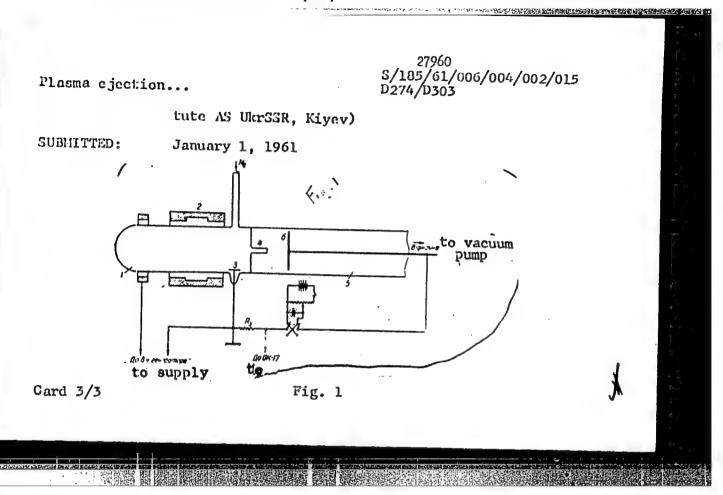
and ion streams on the voltage U of the electrode 6. From the oscillograms it follows that the potential of the plasma with respect to electrode 6 is a periodic function of time. It turned out that the electron andion pulses were shifted by a half-period. The reason for this could be the changing longitudinal electric polarization of the plasma. In the above experiments, the voltage of electrode 6 was given with respect of plate 4, made of metal. Further experiments were carried out, in which the voltage was given with respect to reference electrode 3, made of glass. In this case, the maximum ion-stream density at the outflow from the channel was well above 100 amp/cm². The velocity of the plasma was estimated by the magnitude of the time-shift of the pulse fronts related to the displacements of electrode 6. The velocity was found to be approximately 8·10⁶ cm/sec., corresponding to a proton energy of nearly 30 eV. The charge concentration was found to be approximately 1·10¹⁴ cm⁻³. There are 6 figures and 6 references: 5 Soviet-bloc and 1 non-Soviet-bloc.

ASSOCIATION:

Instytut fizyky AN USSR, m. Kyyiv (Physics Insti-

Card 2/3

"APPROVED FOR RELEASE: 03/13/2001 CIA-RDP86-00513R000513920017-3



20668

S/057/61/031/001/013/017 B104/B204

26.2021

Gabovich, M. D., Pasechnik, L. L., and Romanyuk, L. I.

TITLE:

AUTHORS:

The boundary of a penetrating plasma and plasma focusing

PERIODICAL: Zhurnal tekhnicheskoy fiziki, v. 31, no. 1, 1961, 87-93

TEXT: The authors describe a probing method for determining the boundary of a penetrating plasma. The experimental arrangement shown in Fig. 1 consists of ε pulsed ion source with electron oscillations in a magnetic field. The discharge current attains 40 a, the ion pulses have a rectangular shape, the pulse repetition frequency is 50 cps, and the magnetic field strength is about 300 cersteds. The discharges were produced in hydrogen at a pressure of 5°10" mm Hg. The plasma coming from the source passes through an opening in an electrode (9), and reaches a lens consisting of two cylinders (10) and (11) (inner diameter of the cylinder: 120 mm; L = 120 mm; distance ΔL: 20 mm). Electrode (11) has a negative potential of U = 50 ky relative to electrode (10). A beam catcher

prevents secondary electron emission from electrode (11). Probes (7) and (8) could be shifted. The signal coming from the probes was amplified Card 1/7

- 20668

The boundary of a penetrating

S/057/61/031/001/013/017 B104/B204

and fed into a peak generator. The output signal of this peak generator was conveyed to a recorder, whereby the spatial distribution of the probe current could be recorded. From the axial and radial distributions of the plasma parameters near the opening, which are shown in Figs. 4 and 5. it follows that an increase of the negative potential of electrode (11) up to U = 30 kv produces no effect upon the distribution of the plasma parameters. At a greater distance from the opening, determination of the plasma parameters is more difficult. The authors confined themselves to determining the plasma boundary, and, for this purpose, they applied a potential of 100 v to the probe relative to electrodes (5) and (9); the probe current was automatically recorded. In this way, a plasma boundary could be clearly determined. This boundary is at a distance of about 10-15 mm from the opening and manifests itself in a change in the drop of the probe current. Up to approximately 10 mm, the probe current drops exponentially; at larger distances a greater drop occurs (Fig. 6). In this way, it is possible to determine the plasma boundaries for various conditions. As may be seen from a close study of the plasma boundaries, the shape and position of the plasma boundary change with a change in U,, which is equal to a change in the focusing properties of the system. Card 2/7

The boundary of a penetrating ...

\$/057/61/031/001/013/017 B104/B204

If the plasma boundaries are simulated with metal electrodes of corresponding configuration, it is possible, conditions being suitable, to construct the ion trajectories (Fig. 9). From this figure it may be seen that by increasing the potential and extending the plasma boundary, the ion current focused in the beam catcher may be increased. Fig. 10 graphically represents the experimental dependence of the ion current on the potential U. There are 12 figures, 1 table, and 7 references: 4 Soviet-bloc and 2 non-Soviet-bloc.

ASSOCIATION: Institut fiziki AN USSR Kiyev

(Institute of Physics AS UkrSSR, Kiyev)

SUBMITTED:

June 1, 1960

Card 3/7

"APPROVED FOR RELEASE: 03/13/2001

CIA-RDP86-00513R000513920017-3

24, 2120 (1049,1482,1502) 10.8000

2092lı \$/057/61/031/003/008/019 B125/B202

26.2021

AUTHORS:

Gabovich, M. D. and Romanyuk, L. I.

TITLE:

Effect of a magnetic field on the shape of the boundary of a penetrating plasma and on plasma focusing

PERIODICAL:

Zhurnal tekhnicheskoy fiziki, v. 31, no. 3, 1961, 315-320

TEXT: The authors demonstrate that a magnetic field may considerably influence the shape of the boundary of a penetrating plasma. According to the configuration of the boundary concerned the magnetic field may improve or impair plasma focusing. The development of new methods of controlling the shape of the plasma boundary is of concern. The apparatus used for these experiments has been described already by M. D. Gabovich, L. L. Pasechnik, L. I. Romanyuk, (ZhTF, 31, 87, 1961). It is illustrated once again in Fig. 1. Like in earlier studies the authors used a pulsed ion source with a duration of pulses of 100 microseconds and with a frequency of 50 pulses/sec. In this case the plasma penetrated into the plasma lens consisting of electrodes 6 and 7 through a hole in electrode 4. The plasma boundary was determined by two probes 2 and 3. Fig. 2

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20924

Effect of a magnetic field on the ...

S/057/61/031/003/008/019 B125/B202

shows the dependence of the amperage Ly measured in the beam catcher on the magnetic field strength produced by the solenoid. With an intensification of the magnetic field I decreases already with very weak magnetic fields. Fielis with some dozens of oersteds are already sufficient for a considerable reduction of the ion current focused into the beam catcher. The configuration of the boundary of the penetrating plasma is changed as a result of its contraction and may impair the focusing properties of the system. Also a weak magnetic field may disturb plasma focusing, however, at least two cases exist where the magnetic field improves the focusing of the ions: 1) Focusing with lacking magnetic field under exclusive action of an electric field. 2) If the magnetic field in the discharge chamber of the source considerably penetrates into the region of the plasma to be studied. In the last chapter the author describes a ring-shaped plasma source. The following problem is dealt with: Let us replace the sole opening with its center on the axis of the source by several openings which lie on a concentric circle. Is the intensity of the plasma near the axis of the lens weakened and is the concentration thus distributed over the plasma surface? In what manner is the quality

Card 2/4

20924 \$/057/61/031/003/008/019 B125/B202

Effect of a magnetic field on the ...

replaced by 6 openings lying on a circle with a radius of 19mm; the ion source, however, remained the same. In the case of the holes circularly arranged the density of the ion current was considerably lower than with one central opening (in the cases studied here $j_p = 65 \text{ ma/cm}^2$ and $j_p = 440 \text{ ma/cm}^2$). In the case of circularly arranged holes more than 70% of the total ion flux could be focused into the beam catcher. With the concentrically circularly arranged holes the configuration of the plasma boundary is much more concave than in the case of a single central opening. Also in this case the magnetic field impairs the focusing of the plasma since the plasma is contracted and a projection is formed on the concave boundary of the plasma. The authors conclude that the best results can be obtained by passing the plasma through openings which are at a certain distance from the axis of the source. In some cases such systems are less sensitive to the effect of magnetic fields. There are 12 figures and

of plasma focusing affected? For this purpose the central opening was

ASSOCIATION: Institut fiziki AN USSR, Kiyev (Institute of Physics AS UkrSSR, Kiyev)

Card 3/4

3 Soviet-blog references

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23722 \$/057/61/031/006/006/019 B116/B203

26.1410

Gaberiah. M. D. and Mitropan, I. M.

TITLE-

Chaegration of hydroxagnetic oscillations in the plasma of

on electrode less pulsed discharge

FERIODICAL: Zhurnal tekhnicheekcy fiziki. v. 51. no. 6, 1961, 676-679

TEXT: The radial cacillations of a plasma connected with the annular furrant formed in an electrode-less pulsed discharge were investigated in various papers, exectably by G. B. F. Hiblett, T. S. Green (Ref. 3: Proc. Phys. Soc. 74, 737, 1939). Here, the results of some experiments are presented. The electrode-less pulsed discharge was excited in a quartz or glass tube of 65 nm dismeter, surrounded by two parallel-connected copper coils V4 and V9 with one winding each (Fig. 1). A battery of 10-mf capacitors thanked to 20 km was discharged into the circuit consisting of the said coils, the lead sizes, and the discharger. The inductance of the current circuit was 0.1 miscalery so that the current coolilation period was about 6 mass. The was was previously ionized by a high-frequency inscharge. The coil K placed in the quartz tube served as a magnetic

Cart 1/5

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Observation of bydromagnetic

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profes and permitted the field strength on the axis of the system to be resoured. The Bourskiy belos F, and F inside plass tubes permitted an

observation and measurement of the annular surrents in the gas in two parts of the tube for the well and on the exial, as well as an estimation of radial objects of the turnent ring. These beits could be replaced by one belt which measured the total annular ourrent in the gas. Another bolt served for measuring the current in the apper winding. The circuits of the magnetic perha and of the bette contained integrating RC elements. The latter were throom to as to observe, on the oscilloscope screen, the amberages and mushetto fields, and not their derivatives. The resulting tabilinarams show that the plasma formed in the discrarge (discharge in hydrogen at a pressure P=0.12 mm Hg) influences the strength of the magnetic field considerably. At first, the magnetic field easily enters the plasma. An increasing phase shift occurs between the two abovementioned quantilies with transition from the second to the third halfcycle. At the beginning of the third half-cycle, the field on the plasma axis has a firection opposite to the outer field. In the first quarter of the third half-cycle, the strength of this phase-smifted field increases strongly. This may be assumed to be due to the formation of a current Card 2/5

"APPROVED FOR RELEASE: 03/13/2001

CIA-RDP86-00513R000513920017-3

23722

Observation of hydromagnetic ...

S/057/61/031/006/006/019 B116/B203

ring, its motion toward the center, and the pinch of the magnetic flux "frozen" in the plasma. In such a pinch of the magnetic flux, the formation of radial oscillations of the plasma connected with the current ring may be expected. This is confirmed by data obtained with belts P_1 and P_2 . The oscillograms obtained with the cuter belt P_1 show that the increase in strength of the magnetic field is preceded by the formation of a plasma ring of some dozen ka near the wall. The oscillograms recorded by P_1 and P_2 show a shift of the plasma ring formed on the wall toward the

axis. The radial oscillations of the plasma are observed after deformation of the current ring. This is shown by the oscillation of currents and the fluctuations in strength of the magnetic field observed at the beginning of the increase of the "frozen" magnetic field; at that time, the plasma layer connected with the current ring is near the middle of the tube radius. After a few oscillations, the plasma ring may shift toward the tube center because of the weakening of the magnetic field. The plasma ring is shifted toward the tube center when the current in the copper windings approaches its maximum, whereupon the current ring decomposes, the conductivity of the plasma decreases, and the outer field.

Card 3/5

23722

Observation of hydromagnetic ...

S/057/61/031/006/006/019 B116/B203

enters the whole plasma. Here, a sudden change of the field (negative to positive) is observed. It is a characteristic feature that the oscillations of the plasma ring occur during the pinch of the magnetic field opposite to the outer field, i.e., of the field, the lines of which are connected with the inner currents in the plasma, and not with the current in the cuter winding. It is shown that the observed period of oscillations coincides, as to the order of magnitude, with the period expected according to the approximation of Ref. 3. It is pointed out that in the case of a heavy gas of the same pressure, the oscillation period is larger, which is also confirmed by the oscillograms obtained (hydrogen was replaced by krypton). The increase of the negative field observed in the discharge in hydrogen and the oscillations of the plasma ring were not observed in the discharge in krypton, as had been expected. The present paper was read at the Vtoroye soveshchaniye po teoreticheskoy i prikladnoy magnitnoy gidrodinamike (Second Conference on Theoretical and Applied Magnetohydrodynamics) in Riga on June 30, 1960. There are 7 figures and 5 references: 4 Soviet-bloc and 1 non-Soviet-bloc.

Card 4/5

23722

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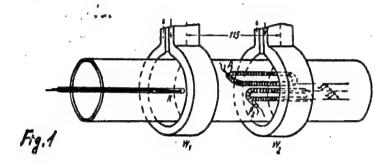
Observation of hydromagnetic ...

ASSOCIATION: Institut fiziki AN USSR Kiyev (Institute of Physics of the

AS UkrSSR Kiyev)

SUBMITTED:

July 25, 1960



Card 5/5

27166 \$/057/61/031/009/006/019 B104/B102

26.2311

AUTHORS:

Gabovich, M. D., Pasechnik, L. L., and Lozovaya, Ye. A.

TITLE:

Discharge of a plasma with high concentration of charged

particles into a vacuum

PERIODICAL:

Zhurnal tekhnicheskoy fiziki, v. 31, no. 9, 1961, 1049-1056

TEXT: The authors studied, by a probing method, the spatial distribution of the parameters of a hydrogen plasma with high concentration of charged particles (about 10¹⁵ cm⁻³). The plasma was produced by a pulsed discharge, the amplitude of the discharge current being about 50 ka. The oscillation period was about 25 µsec, the battery of condensers had 90 µf capacity, and was charged to 3 kv. The most important parts of the experimental arrangement were the plasma source (discharge space with 3 electrodes) and the empty space beyond the hole in the lowest electrode (cf. Fig. 1), where one or two probes could be shifted. All measurements were made at a hydrogen pressure of 5.6°10 mm Hg in the source, and about 10⁻⁵ mm Hg outside the source. In all cases the oscillograms of

Card 1/6

27166 \$/057/61/031/009/006/019 B104/B102

Discharge of a plasma with high. ...

the probe current were recorded together with those of the discharge current. Some peculiarities turned up in the transition from ionic to electronic current; in particular, a strong modulation of the electronic current took place. Such a modulation was observed when the probe exhibited a small positive potential with respect to electrode 2 (Fig. 1) Further, it was remarkable that the ionic current peak agreed almost exactly in time with the discharge current peak, while the electronic current peak was considerably shifted against the discharge current. This is explained by the fact that the probe current depends not only on the plasma concentration but also on the potential in the probe space at the given instant. After determining the probe characteristics, the authors determined the distributions of concentrations of charged particles, of electron gas temperature, and of the space potential. Fig. 9 shows examples of radial distribution of the probe current for distances of the probe from electrode 2 of 5, 10, and 20 mm. Results reveal that the axial distribution of parameters is the same as in plasma with low concentration of charged particles. The temperature gradient is here lower than in plasma with low concentration of charged particles. In the

Card 2/6

27166 \$/057/61/031/009/006/019 B104/B102

Discharge of a plasma with high...

anode cavity, the temperature of the electron gas (about 50,000°K) is lower than in the cathode cavity (130,000-70,000°K). There are 9 figures and 8 references: 6 Soviet and 1 non-Soviet. The reference to the English-language publication reads as follows: The Characteristics of electrical discharges in magnetic fields. Edited by A. Guthrie and R. K. Walkering, N. Y., 1949.

ASSOCIATION: Institut fiziki AN USSR Kiyev (Physics Institute,

AS UkrSSR, Kiyev)

SUBMITTED: August 1, 1960

Fig. 1. Diagram of the experimental arrangement. Legend: 1,2, and 3 are electrodes; 4 is the outlet of the plasma source (3 mm diameter); 5 is the discharger; 7 and 8 are the probes; Co is the capacity for maintaining the probe potential; (A) is an amplifier, (B) an oscilloscope.

Fig. 9. Spatial distribution of the plasma parameters. Legend: (a)

Card 3/6

"APPROVED FOR RELEASE: 03/13/2001 CIA-RDP86-00513R000513920017-3

GABOVICH, M.D. [Habovych, M.D.]; YAZEVA, V.G. [IAzieva, V.H.]

Correlation of low-frequency and high-frequency oscillations induced in a plasma by an electron beam. Ukr. fiz. zhur. 7 no.9:1015-1020 S '62.

(MIRA 15:12)

l. Institut fiziki AN UkrSSR, Kiyev. (Plasma oscillations)

(Electron beams)

1.7216

262312

\$/057/62/032/011/008/014 B104/B102

AUTHORS:

Gabovich, M. D., and Kirichenko, G. S.

TITLE:

The oscillation of ions in the region of the potential minimum and the low frequency oscillations in a gas

discharge

PERIODICAL:

Zhurnal tekhnicheskoy fiziki, v. 32, no. 11, 1962, 1376-1381

TEXT: The oscillation properties of a vacuum diode having very low cesium vapor pressure and those of a mercury vapor diode, are investigated using an experimental arrangement as shown in Fig. 1. The positive ions of cesium produced by thermal ionization move to and foo between the cathode and a point of reversal which depends on the anode potential. If the potential minimum (eV_m \gg kT) is deep, the motion of the ion is determined

not by its thermal velocity but by the electric field. The frequency of oscillations is calculated for pressure so low that the space charge of the cesium ions does not affect the motion of the ions. Application of Langmuir's relations (Phys. Rev., 21, 419, 1923) to the dimensions of the potential well leads to the following formula for the frequency of Card 1/4

The oscillation of ions in the ...

S/057/62/032/011/008/014 B104/B102

the flight ion oscillations:

$$f = \frac{\pi^{9/4}}{2^{1/4}} \frac{1}{\xi_- + \xi_+} \frac{m^{1/4} e^{1/4}}{M^{1/4} (kT)^{1/4}} j^{4/4}.$$

Here γ , $\frac{1}{1+}$ and $\frac{1}{1-}$ are dimensionless parameters tabulated by Langmuir, m is the electron mass, M the ion mass, and j the density of the electron current. The frequency is found to be proportional to $j^{1/2}$ (Fig. 2). The oscillations are independent of the parameters of the external current circuit, are sinusoidal, and have constant frequency and amplitude modulation. The excitation of the oscillations is explained as follows: In the region of the potential well an oscillating electric field exists which modulates the ion velocity and consequently the ion current. This causes the depth of the potential well to pulsate and the anode current to oscillate. If the cathode is incandescent, have to traverse a double layer the positive ions in the plasma of a gas discharge to reach it. This double layer lies near the cathode. Inside the double layer there exists a potential minimum, the depth and dimensions of which can be Card 2/4

"APPROVED FOR RELEASE: 03/13/2001 CIA-RDP86-00513R000513920017-3

The oscillation of ions in the ...

S/057/62/032/011/008/014 B104/B102

determined from the anode current and the saturation current of the cathode in the same way as for the vacuum diode provided that the positive space charge is negligible. The flight-oscillations of the positive ions that occur under these conditions are the main cause of low frequency oscillations appearing in discharges with incandescent cathodes. There are 6 figures.

ASSOCIATION: Institut fiziki AN USSR, Kiyev (Institute of Physics

AS UkrSSR, Kiyev)

SUBMITTED: October 5, 1961

Fig. 1. Block scheme. Legend: (1) commutator of the filament current, (2) electronic commutator, (3) superheterodyne, (4) oscillograph, (5) 9HO-1 (ENO-1) oscillograph, (6) broad band amplifier, (7) cathode follower, (8) OK-17 (OK-17) oscillograph, (9) amplifier.

Fig. 2. Dependence of the frequency on \sqrt{j} for different filament currents. Legend: (1) 13.0 a, (2) 12.2 a, (3) 12.0 a, (4) 11.25 a.

Card 3/4

38858

s/056/62/042/006/010/047 B104/B102

26.2531 26.2312

Gabovich, M. D. Kirichenko, G. S.

17

TITLE:

AUTHORS :

The excitation of oscillations on the passage of a beam of slow ions through a plasma

PERIODICAL:

Zhurnal eksperimental noy i teoreticheskoy fiziki, v. 42, no. 6, 1962, 1478 - 1480

TEXT: A discharge was produced in cesium vapor (N 10⁻³ mm Hg) between the cathode and the anode (Fig. 1). The plasma of this discharge penetrated the side branch with probes, 3_2 , 3_1 , and 3. The probes 3_1 and 3_2 are placed at distances of 9 and 15 mm from the probe 3. A positive potential relative to the plasma was applied to 3. At high temperatures of 3, cesium ions were produced by surface ionization. Owing to the potential difference these ions fell into the plasma as a straight line beam. In agreement with the theory, low frequency oscillations were excited on the passage of slow ions through the plasma. The low frequency oscillations made their appearance at those temperatures of 3 at which desium ions were produced. The frequency of the oscillations lies in the region of the

Card 1/2

ACCESSION NR: AT4025320

8/0000/63/000/000/0283/0291

AUTHORS: Gabovich, M. D.; Kirichenko, G. S.; Koydan, V. S.

TITLE: Excitation of plasma oscillations by an ion beam, and the possibility of determining the electron temperature

SOURCE: Diagnostika plazmy* (Plasma diagnostics); sb. statey. Moscow, Gosatomizdat, 1963, 283-291

TOPIC TAGS: plasma oscillation, ion beam, plasma ion oscillation, plasma electron temperature, plasma interaction, drift, standing wave

ABSTRACT: Continuing their earlier investigations ("Zh. eksperim. i teor. fiz." v. 42, 1478, 1962; Ukr. fiz. zh., in press), the authors describe apparatus aimed at checking the influence of electron drift in a direction opposite to the ion beam on the stability of the oscillations produced when an ion beam passes through a plas-

Card 1/3

CIA-RDP86-00513R000513920017-3

ACCESSION NR: AT4025320

The apparatus and its operation are briefly described. It is reported that, unlike the earlier experiments, oscillations with noticeable amplitude were excited also in the absence of drift cur-The properties of these oscillations are described briefly. In the presence of backward drift, a new type of more intense oscillation with a rather narrow frequency spectrum was also observed. It is concluded that the backward electron drift leads to establishment of a standing wave, to a considerable increase in the oscillations, and to a narrowing down of the frequency range. The ion threshold energy at which the excitation of these oscillations terminates is proportional to the electron temperature. This is in qualitative agreement with the theory and gives grounds for assuming that a new method will be developed for determining electron temperature. It is proposed in the future to broaden the range of electron temperatures of the investigated plasmas and also to carry out a rigorous quantitative determination of the threshold energy. Orig. art. has: 7 figures.

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1."

"APPROVED FOR RELEASE: 03/13/2001 CIA-RDP86-00513R000513920017-3

ACCESSION NR: AT4025320		
ASSOCIATION: None		
SUBMITTED: 190ct63	DATE ACQ: 16Apr64	ENCL: 00
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"APPROVED FOR RELEASE: 03/13/2001

CIA-RDP86-00513R000513920017-3

AID Nr. 995-11 21 June

PLASMA ION SOURCES AND ION-BEAM FORMATION (USSR)

Gabovich, M. D. Pribory i tekhnika eksperimenta, no. 2, Mar-Apr 1963, 5-19.

S/120/63/000/002/001/041

A detailed review of the state of research and development of 1) ion extraction from plasma ion sources and 2) primary formation of ion beams is given, based on 59 articles (27 Soviet) published up to October 1962. The conditions of three methods of ion extraction are described: 1) the emissive surface of the plasma is flat, within the ion source, and close to the outlet aperture; 2) the emissive surface is concave, within the ion source, and much larger than the cross section of the outlet channel carrying the focused ion beam; and 3) the plasma extends beyond the ion source and forms a well-developed emissive surface ahead of the focusing lens. Each type is subjected to a broad analysis and interpretation of its potentialities.

Card 1/1

S/185/63/008/001/008/024 D234/D308

AUTHORS:

Habovich, M. D., Lozova, O. O. and Romanyuk, L. I.

TITLE:

Possibility of location of the boundary of penetrating

plasma by a beam of charged particles

PERIODICAL:

Ukrayins'kyy fizychnyy zhurnal, v. 8, no. 1, 1963,

57-59

TEXT: If a beam of electrons passing through plasma and falling on a fluorescent screen is displaced away from the ion source, the bright spot on the screen will also be displaced in the same direction until the beam reaches the plasma boundary, and then in the opposite direction owing to the reflection of the beam at the boundary. By varying the inclination of the beam one can determine the position and the shape of the boundary. The authors describe an experimental installation which they used for checking this method. Data agree well with those obtained by the probe method if the potential is not too high. The error at high potentials is explained by the fact that the boundary becomes convex, and use of Card 1/2

"APPROVED FOR RELEASE: 03/13/2001 CIA-RDP86-00513R000513920017-3

Possibility of location ...

S/185/63/008/001/008/024 D234/D308

heavy negative ions instead of electrons is suggested in this case. There are 3 figures.

ASSOCIATION:

Instytut fizyky AN URSR (Institute of Physics of the AS UkrSSR), Kiev

SUBMITTED:

August 3, 1962

Card 2/2

S/109/63/008/003/022/027 D271/D308

AUTHORS:

Gabovich, M. D., and Kirichenko, G. S.

TITLE:

Microwave generation by electrons oscillating in the potential well formed by positive space charge

cnarg

PERIODICAL:

Radiotekhnika i elektronika, v. 8, no. 3, 1963, 520 - 522

TEXT: A microwave generator is discussed, which consists of a diode with an ion current limited by space charge; the frequency does not depend on the external circuit but depends only on the collector potential. In a hot-cathode diode with cesium vapor, conditions can be created in which a potential well is formed near the cathode, and cesium ions oscillate in this well. The frequency of oscillations depends on emitter temperature, electrode spacing, potential difference, and the mass of oscillating ions. Frequency-voltage dependence is shown in a graph.

Card 1/2

Microwave generation...

S/109/63/008/003/022/027 D271/D308

A diode of this type can be used for microwave generation if the collector is made negative and cesium ions from the cathode are attracted towards it. Owing to the positive space charge limiting the current, a potential well is formed in which electrons may oscillate. The frequency is calculated as in the case of ion oscillation, except that electron mass is substituted for ion mass. The temperature of the emitter must be such as to produce both ions and electrons. In an experimental diode with indirectly heated tantalum emitter and cesium vapor, oscillations were detected in the range of 150 - 1000 Mc/s in a narrow temperature range around 1900 K. Frequency was found to be in good agreement with the analytical value. Wide frequency variations are possible. There is 1 figure.

SUBMITTED:

August 13, 1962

Card 2/2

"APPROVED FOR RELEASE: 03/13/2001 CIA-RDP86-00513R000513920017-3

ACCESSION NR: AL	IJP(C)/SSD Pz-4			2
			/63/008/006/0624/0627	
AUTHOR: Gaboure	h, M. D., Kytry*chen	ko, G, S,	·	i
1			of plasma oscilations by	
SOURCE: Ukrains	kyy fizychnyy zhurna	1, v. 8, no. 6, 19	63, 624-627	
TOPIC TAGS: plas	MA Osoillation ton	hoom American	. ion beam-plasma inter- el, plasma excitation.	G Valence
ABSTRACT: The pu	roose of this investi	t-att	ify the theory of stability ry predicts the criterion	
	•	$+\left(\frac{\omega_{+R}}{\omega_{+RA}}\right)^{5/2}$		
there V is the vel	locity of the beam ic	ons, C _e the thermal	speed of the electron S.	
ard 1/42			•	

L .18247-63....

ACCESSION NR: AP3002115

An experimental setup was devised by which it was possible to control the ion energy. Oscillations were measured by means of a probe up to a frequency of 4.35 Mes at various values of delta V, the difference of potential between the emitter and the anode. At delta V value less than 3.5 Volt ions can not reach the probe, and only a noise was registered. At delta V between 4.5 and 13.5 Volt the ion beam excited plasma, so that definite oscillation peaks were observed on the spectrum. As the beam energy increased the oscillation amplitude decreased, and at a voltage exceeding 13.5 Volt it was impossible to observe oscillations because of the noise background. The oscillation amplitude and frequency are shown on fig. 3 of enclosure Ol as functions of ion energy E. The E values were calculated from delta V, the space potential and the contact potential difference between the probe which is coated with a Cs film and the emitter which has no such coating. The measured temperature T sub e reached about 20,000 K. The concentration of beam ions and plasma ions was about equal. The ion energy E was then at its highest value of about 6kT sub e. And when this value was exceeded the system became stable. The experiments substantiated the theory. The orig. art. has: 1 formula, and 3 figures.

cord 2/1/2 Physics INSTITUTE OF UKR SSR. RAD. of S.

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WM5011014 BOOK EXPLOITATION UR/

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Clabovich, Mark Davidovich

Plasma ion sources (Plazmennyye istochniki ionov), Kiev, Naukova dumka, 1964, 220 p. illus., biblio. (At head of title: Akademiya nauk Ukrainskoy SSR. Institut fiziki) 2,800 copies printed.

TOPIC TAGS: plasma, ion, plasma oscillation, ionized plasma, plasma source, ion source, hydrogen ion, ion beam, gas discharge

FURFOSE AND COVERAGE: This book is a survey of literature on gas discharge ion sources. Different types of plasma ion sources are examined: sources of atomic and molecular hydrogen ions and multiply charged ions, electrons and other particles. This study is also concerned with ion extraction and the primary shaping of ion beams. The yield of quasi-neutral ion flows in a vacuum and the role ofoscillation properties of plasma are also considered. The book is intended for scientific, engineering, and technical personnel and students.

TABLE OF CONTENTS (abridged):

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Ch. II.		trum and charge com	position of beam	s derived from p	olasma
Ch. III	Ion extracti	on and the primary	shaping of ion b	eams. Beam prope	gation
	Penetration of leaving the p	f plasma from the i lasma of the source	177		
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ACCESSION NR: AP4020578

S/0057/64/034/003/0488/C495

AUTHOR: Gabovich, M.D.; Romanyuk, L.I.; Lozovaya, Ye.A.

TITLE: Escape of plasma from an oscillating electron source into vacuum in the presence of a magnetic field

SOURCE: Zhurnal tekhnicheskoy fiziki, v.34, no.3, 1964, 488-495

DDPIC TAGS: plasma, plasma source, oscillating electron plasma source, plasma in magnetic field, probe measurements, thermal probes, plasma escape

ABSTRACT: The escape of a helium plasma from an oscillating electron source into vacuum was investigated experimentally in the presence of a magnetic field. The source employed a 6-mm diameter indirectly heating cathode on the axis of a 3-cm diameter cylindrical anode. The reflecting electrode was located 6 cm from the cathode, was kept at cathode potential, and had a 3-mm diameter opening for plasma escape. The glass vacuum chamber was about 12 cm in diameter and 27 cm long. Gas pressures of 2×10^{-2} and 2×10^{-4} mm Hg were maintained in the source and the vacuum chamber respectively. Anode potentials from 150 to 200 V were employed with discharge currents from 1.0 to 1.5 A. The source and vacuum chamber were located in a

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uniform longitudinal magnetic field of 1000 Oe or less. The escaping plasma was investigated with probes of various types. In spite of the strong magnetic field, the ion current in the escaping plasma was not confined to the axis of the chamber but extended several centimeters from the axis. The ion current was due mostly to ordered motion, the current due to chaotic motion being very small. Most of the ions had energies roughly equal to the cathode drop in the discharge. There was a small admixture of lower energy ions. The distribution of electrical potential in the escaping plasma was determined with the aid of two types of thermal probe. At a fixed distance from the source the potential, as a function of the radius, showed a minimum on the axis of the chamber and a maximum some millimeters off the axis. On the axis the potential (with respect to the cathode and reflector) was large and positive near the source and fell rapidly to zero within a few centimeters. At the axis of the chamber an insulated probe assumed a large negative potential of several tens of volts. This potential increased in absolute value (became more negative) as the distance from the scurce was increased. When the probe was moved off the axis, the potential first fell rapidly to zero and then became positive. This behavior is interpreted as indicating the presence of a narrow beam of fast electrons produced by interaction of the electron current with the plasma within the source. Orig.art. has: 3 formulas and 7 figures.

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3/0057/64/034/006/0993/0997

ACCESSION NR: AP4040299

AUTHOR: Gabovich, M.D.; Mitropan, I.M.

TITLE: Interaction of plasma streams moving in opposite directions along the axis of an induction pinch

SOURCE: Zhurnal tekhnicheskoy fiziki, v.34, no.6, 1964, 993-997

TOPIC TAGS: plasma, plasma containment, plasma stability, hydrogen plasma

ABSTRACT: The behavior of an induction pinch in a hydrogen plasma was investigated. The pinch was produced in a 6.4-cm-diameter glass tube by discharge of a 10-micro-farad capacitor bank, charged to about 25 kV, through two copper loops encircling the tube and located 8.4 cm apart. The period of this system was somewhat greater than 6 microsec. The magnetic field at the axis of the tube was measured with a movable probe. When the hydrogen pressure was 0.1 mm Hg, the phenomena observed were the same as previously described by the present authors (ZhTF 32,1371,1962). At a pressure of 0.06 mm Hg, the oscillogram from the magnetic probe was the same as before, when the probe was located beneath one of the copper loops, but it altered as the probe was moved toward the central plane between the two loops. In this

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ACCESSION NR: AP4040299

central region the magnetic field showed two maxima and fell to zero between them. This behavior is ascribed to an instability of the type discussed by W.E.Nixon, W. F.Cummings, F.H.Coensgen, and A.E.Sherman (Phys.Rev.119,1457,1960) due to the intermingling of two streams of plasma flowing in opposite directions toward the central plane from the regions of high magnetic field beneath the loops. The experiment was repeated with a constriction 1 cm in diameter and 3.5 cm long in the discharge tube and a specially constructed differential flux meter entirely outside the tube. Similar results were obtained. When one end of the constriction was closed with a glass stopper, preventing flow of plasma toward the central plane from one direction, the diamagnetic effect disappeared. Orig.art.has: 5 figures.

ASSOCIATION: none

SUBMITTED: 25Jun63

ATD PRESS: 3084

ENCL: 00

SUB CODE: ME, EM

NR REF SOV: 002

OTHER: 004

2/2

L 14480-65 ENT(1)/ENG(k)/EPA(sp)-2/EPA(w)-2/EEC(t)/T/EEC(b)-2/ENA(m)-2 Pz-6/Po-4/Pab-10/Pi-4 IJP(c)/AFWL/SSD(t)/AEDC(b)/SSD/ASD(a)-5/BSD/ASD(f)-2/ASD(p)-3/AFETR/RAEM(a)/ESD(gs)/ESD(t) AT ACCESSION NR: AP4047932 S/0056/64/047/004/1594/1595

AUTHOR: Gabovich, M. D.; Kirichenko, G. S.

TITLE: The threshold energy of ions in two-beam ionic instability

SOURCE: Zhurnal eksperimental noy teoreticheskoy fiziki, v. 47, no. 4, 1964, 1594-1595

TOPIC TAGS: plasma, plasma instability, plasma oscillation, ion beam, electron temperature, plasma diagnostics

ABSTR/CT: It is pointed out that relative motion of an ion beam and plasma or plasma streams can be accompanied by two-beam ionic instability. Oscillations can be excited in passage through the plasma. The present article describes experiments performed in order to determine the threshold energy of ions in two-beam instability. Ion beams with currents on the order of 1 mamp were passed through plasma formed in different gases, and oscillations excited by the beam at different electron temperatures were observed. It was established that the dependence of the amplitude of the oscillations on the energy

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gy of the ions can be represented by a curve which contains a maximum. It was found that in accordance with theoretical predictions, the threshold energy of ions determined from the maxima as a function of the temperature of the ions is proportional to the electron temperature. The frequency of oscillations was determined to be approximately proportional to the velocity of the ions. It is pointed out that the experiments indicate the possibility of using two-beam ionic instability for thermalization of powerful ion beams in plasma with high electron temperature and for plasma diagnostics. Orig. art. has: 1 figure.

ASSOCIATION: none

SUBMITTED: 02Jul64

ENCL: 00

SUB CODE : "NP , HE

NO REF SOV: 005

OTHER: 001

ATD PRESS: 3136.

Card 2/2

L 27599-65 EWT(1)/EPA(sp)-2/EPF(c)/EPA(w)-2/EEC(t)/T/EWA(m)-2 Pz-6/po-4/pab-10/Pr-4/P1-4 IJP(c) WW/AT

ACCESSION NR: AP5003241

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AUTHOR: Gahovich, M.D./ Romanyuk, L.I./ Lozovaya, Ye.A.

72 46 B

TITLE: Formation of a quasineutral beam of accelerated ions in the plasma issuing from an ion source.

SOURCE: Zhurnal takhnicheskoy fiziki, v.35, no.1, 1965, 94-100

TOPIC TAGS: plasmi, ion beam, ion source, ion acceleration

ABSTRACT: This paper reports a continuation of previous work of the authors (ZhTF 34,488,1964) concerning the reflex discharge ion source. The apparatus is similar to that described in the earlier paper, with such modifications as were required for the particular experiments performed. The apparatus was operated under a variety of conditions, the current-voltage characteristics were measured, and particular attention was given to the potential gradient in the plasma beam issuing from the source. The principal conclusion is that the following conditions are requisite for obtaining ions with energies corresponding to the cathode drop: the issuing plasma must contain an intense beam of primary electrons with appropriate velocity distribution; the plasma must issue from the chamber into a region of suf-

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ACCESSION NR: AP5(103241

ficiently high vacuum; there must be not positively charged electrode outside the discharge chamber that could remove electrons from the issuing plasma. Orig.art. has: 6 figures.

ASSOCIATION: Institut fiziki AN UkrSSR, Kiev (Institute of Physics, AN UkrSSR)

SUBMITTED: 24Feb64

ENCL: 00

SUB CODE: ME.NP

NR REF SOV: 005

OTHER: 003

Card2/2

L 41606-66 EWT(1) IJF(c) AT

ACC NR: AF6018796 . SOURCE CODE: UR/0056/66/050/005/1183/1186

AUTHOR: Gabovich, M. D.; Kirichenko, G. S.

7

ORG: Institute of Physics, Academy of Sciences, Ukrainian SSR (Institut fiziki Akademii nauk Ukrainskoy SSR)

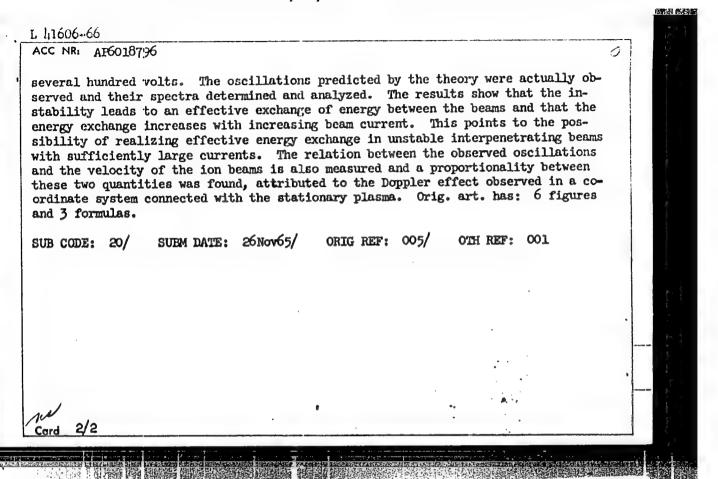
TITLE: Two-stream instability in a system of interacting ion beams

SOURCE: Zh eksper i teor fiz, v. 50, no. 5, 1966, 1183-1186

TOPIC TAGS: plasma instability, ion beam, plasma beam interaction, plasma electron temperature, plasma oscillation, Doppler effect

ABSTRACT: This is a continuation of earlier work by the authors (ZhETF v. 47, 1594, 1964 and elsewhere) dealing with the interaction of an ion beam with a plasma. These investigations have verified the basic theory of two-stream ion instability and have demonstrated the possibility of thermalization of an intense ion beam in a plasma characterized by a high electron temperature. The present study is devoted to instability of interpenetrating potassium ion beams with energies up to 4 kev in a plasma formed by the ionization of a gas (krypton or neon) at a pressure 3 x 10⁻⁶ - 10⁻⁴ mm Hg by these fast ions. It is shown that two-stream ion instability can arise in such a system, which can be regarded as consisting of two ion beams moving in the same direction but with different velocities, if the energy difference in the beams is smaller than some threshold value. For example, for ion beams with energies of the order of several kev, with electron temperature of 1 ev, the threshold may be

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L 06310-67 ENT(1)/ENT(m)/ENP(t)/ETI IJP(g) AT/JD/JG/ED ACC NR: AT6020434 (N) SOURCE CODE: UR/0000/65/000/00044/0051	1
AUTHOR: Gabovich, M. D.; Kirichenko, G. S.; Koydan, V. S.	
ORG: none	1 (1) K
TITLE: Interaction of ion beams with a plasma	200
SOURCE: AN UkrSSR. Vzaimodeystviye puchkov zaryaznemnyku chasta solomka, 1965, 44-51	
TOPIC TAGS: plasma beam interaction, ion beam, cesium plasma, inert gas, gas density, plasma electron temperature, standing wave	
ABSTRACT: The experimental parameters were chosen to satisfy the instability criterian derived by Vedenov, et al (UFN, 1961, 73, 701) using a cesium ion beam with an energy of several ev. A plasma of 10 cm 3 particle density was produced in inert gas discharges. Beam density was of the same order of magnitude. The amplitude and frequency of oscillations excited by ion beams was studied as a function of the electron temperature, gas density and ion mass. It is shown that the peak amplitude of the frequency spectrum can be explained by the theoretical ion beam energy at which stable opecy spectrum can be explained by the theoretical ion beam energy at which stable opecy spectrum. As magnetic field was increased (in a direction parallel to the beam), ration occurs. As magnetic field was increased (in a direction parallel to the beam), ration occurs increase in noise which made the diagnostic measurement more diffithere was a great increase in noise which made the diagnostic measurement more diffithere was a great increase in show that the excited oscillations have the same	A The Control of the
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r a feedback hat ion beam	scheme which go	enerated a star	iding wave. The exper	"Imental regults show !	
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GABOVICH, M.O.; TSVETSIESKIY, S.V.

Operation of retary centimeus diffusors. Sakh.prom.30 me.3:24-31
Mr '56. (MEA 9:7)

1.Eurskiy sakhsvekletrest (for Gabevich).2.Sakharayy saved imeni
Kuybysheva (for TSvetsinskiy).

(Sugar machinery)

GABOVICH, M.S.

Modernization of rotary diffusers. Sakh.prom. 32 no. 7:8-12 Jy '58. (NIRA 11:8)

l. Upravleniye sakharnoy promyshlennosti Kurskogo sovnarkhoza. (Diffusers)

GABOVICH, H.S.

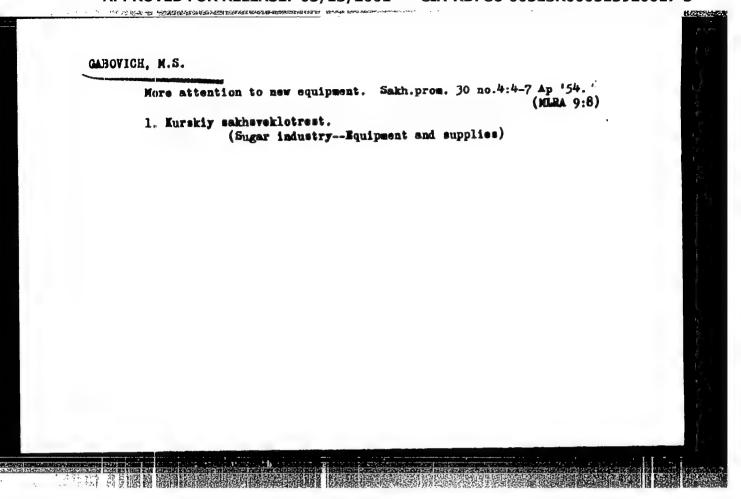
Quality of brass tubes for the surfaces of heat-exchange apparatus. Sakh.prom. 32 no.11:33-34 N '58. (MIRA 11:12)

l.Upravleniye sakharney premyshlennesti Kurskege sovnarkhoza. (Heat exchangers) (Tubes, Cepper)

GABOVICH, M. S.

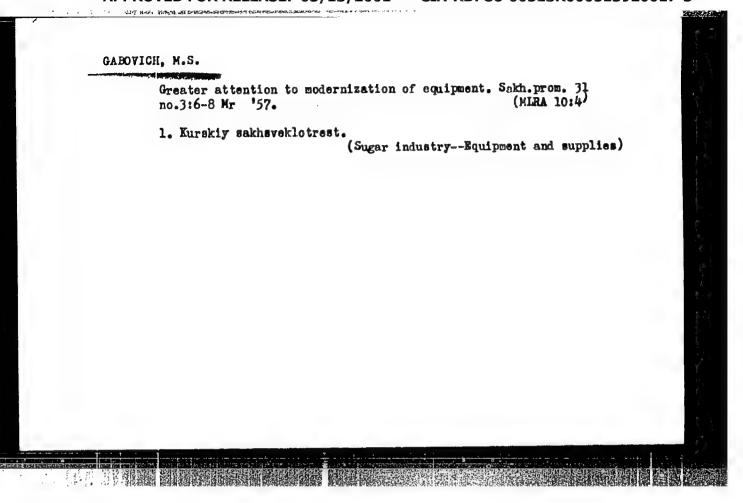
Some defects of new diffusion batteries. Sakh. prom., 26, no. 3, 1952.

Increasing the capacity of heat exchangers. Sakh.prom. 27 no.8:34-35 Ag '53. (NLBA 6:8) 1. Kurskiy sakhsvekletrest. (Sugar machinery)



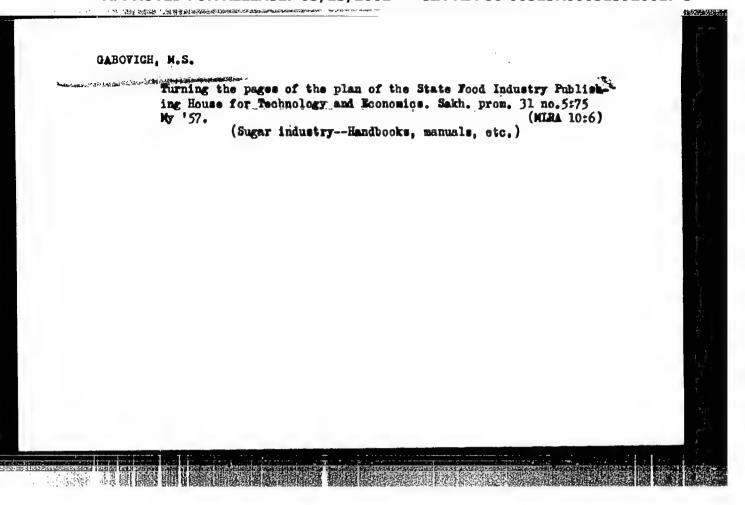
Concerning M.A.Miagkov's article on material and technical supply.
Sakh.prom. 29 no.1:5-6 '55. (MIRA 8:4)

1. Kurskiy sakhsveklotrest.
(Sugar industry—Equipment and supplies)



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CIA-RDP86-00513R000513920017-3



MOSKALENKO, S.I.; GABOVICH, M.S.; BACHINSKIY, Tu.V.; TOMILIN, A.V.;

MEDVEDEV, P.M.; LOMANOVA, M.M.; GOLOVKOV, P.D.; GAYDUKOV, G.I.;

ALEYNIKOV, V.V.; STEMIN, N.D.; MIRONOVA, V.V.; BELAVIMESSVA,

Ye.S.; TSVETSINSKIY, S.V.; NECHEPURNYY, P.; KOBZAR', H.K.;

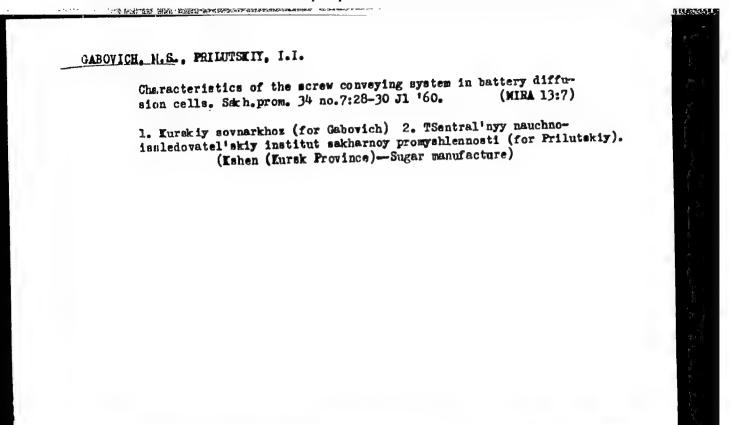
ROZHMOVA, Ye.S.; PHINTMINSKIY, V.N.; GOHDETCHUK, V.K.; SHOKRIGO,

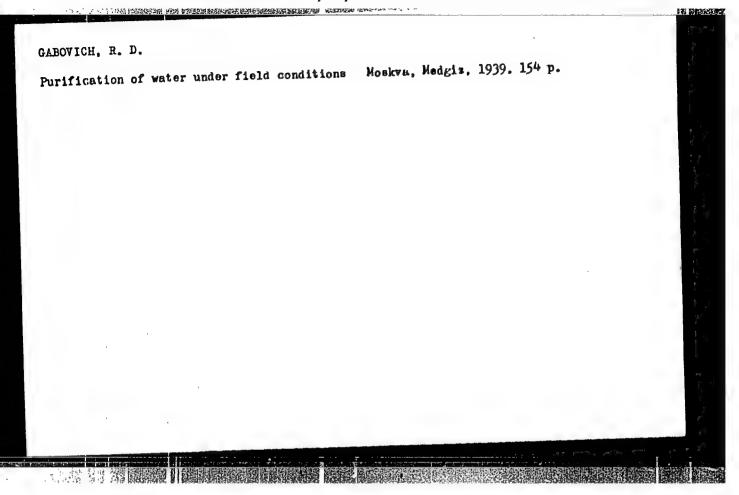
V.F.; KISLTUK, M.

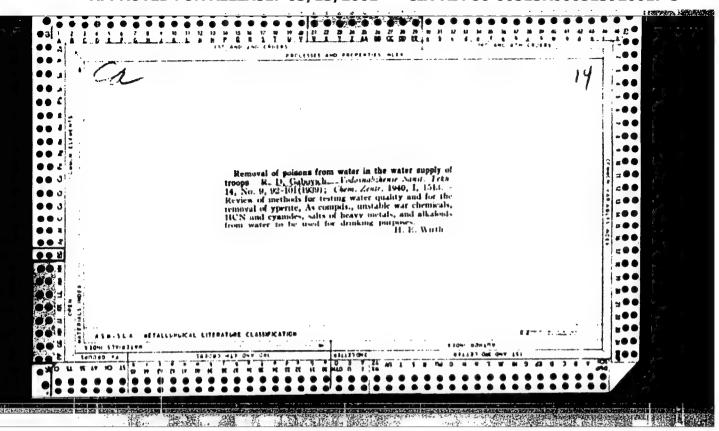
Fifty years in the sugar industry. Sakh.prom. 33 no.2:18

[MIRA 12:3)

(Shtepan, Georgii Viacheslavovich, 1888-)



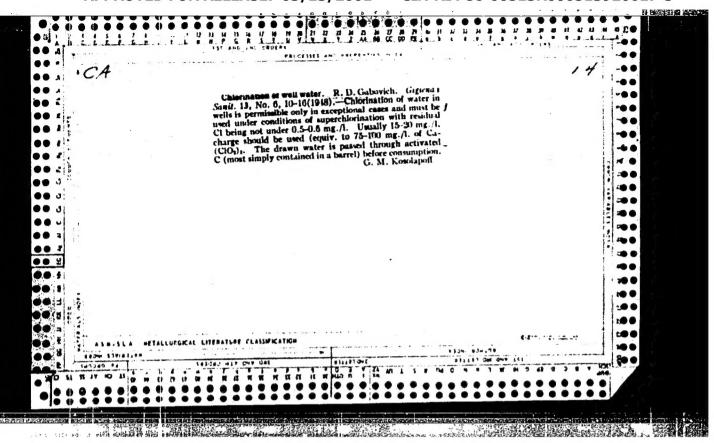




GABOVICH, R. D.

Gabovich, R. D. "Fluorine in the drinking water of the Ukraine and its hygienic significance" (Summary of the paper), Soobshch. o nauch. rabotakh chlenov Vsesoyuz. khim. o-va im. Mendeleyeva, 1948, Issue 3, pp. 9-10.

SO: U-3261, 10 April 53 (Letopis 'Zhurnal 'nykh Statey No.11, 1949)



GABOVICH, R. D.

"Chlorination of Well Water," Gig. 1 San., No.6, 1949

Chair Gen. Hygiene, Kiev Med. Int.

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CIA-RDP86-00513R000513920017-3

USER/Medicine - Water, Supply Jul 19

Medicine - Water Purification

"The Fluorine Content in Drinking Water in the Ukrainian SER," Docent R. D. Gabovich, Chair of Gen Hygiene Kiev Med Inst, 52 pp

"Gig i San" No 7

Fluorine content in Ukrainian water sources varies from a mere trace to over one mg/l. Mass inspection of drinking water has already made it possible to study the relation between fluorine concentration and pitted enamel and caries in teeth in addition to other ailments. This work should be promoted by hygienists.

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